

# Tomographic Image of the Proton from DVCS data

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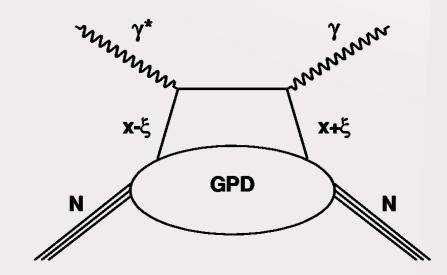
# **Generalized Parton Distributions**

#### A Generalization of usual PDFs

- A description of the nucleon depending on 3 variables : x,  $\xi$ , t
- Accessible through several exclusive processes (DV Compton scattering, DV Meson Production, Double DVCS...)
- Spin-1/2 of the nucleon leads to
   4 GPDs

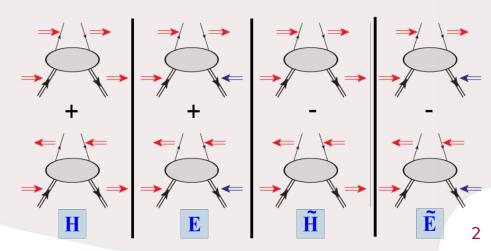
## Deep Virtual Compton Scattering (DVCS)

- Exclusive electroproduction of a photon
- Using the factorization gives access to GPDs
- However x is not observable in DVCS
- Only Compton Form Factors (CFF) are accessible



$$F_{Re}(\xi,t) = \mathcal{P} \int_{-1}^{1} dx \left[ \frac{1}{x-\xi} \mp \frac{1}{x+\xi} \right] F(x,\xi,t),$$

$$F_{Im}(\xi,t) = F(\xi,\xi,t) \mp F(-\xi,\xi,t).$$



M. Guidal et al. Rept. Prog. Phys. 76 (2013) 066202



# **Generalized Parton Distributions**

- Many structure functions are used to understand the structure of the nucleon
- GPDs offer several advantages
  - The link to PDFs is well established
  - Experimental access through DVCS
     A.V. Radyushkin Phys.Lett. B380 (1996) 417-425
- GPDs give access to the proton spin through Ji sum rule

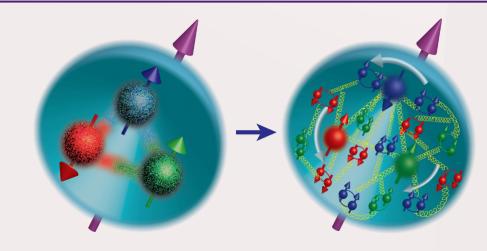
$$J_{q,g} = rac{1}{2} \int_{-1}^{+1} dx \; x [H_{q,g}(x,\xi,t=0) + E_{q,g}(x,\xi,t=0)]$$

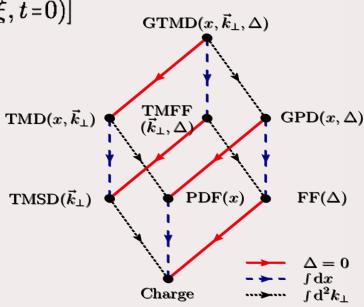
X.D. Ji Phys.Rev.Lett. 78 (1997) 610-613

GPDs give a 3D image of the nucleon

M. Burkardt Phys. Rev. D62 (2000) 071503

- And much more
  - Pressure and shear forces in the D-term
     M. Polyakov Phys.Lett. B555 (2003) 57-62
  - Study of nuclei to understand the EMC effect and shadowing
    - With an access to non-nucleonic degrees of freedom in nuclei







# Disentangling the CFFs

 The 4 complex CFFs intervene as 8 free parameters in the calculation in various observables of DVCS

A. Belitsky et al. Nucl. Phys. B629 (2002) 323-392

- We need many observables to extract all of them
- Possible for one point using HERMES data
  - Includes transversely polarized target and beam charge asymmetries
- JLab data however does not allow for a full extraction of that kind
  - Use the fact that the Im(H) dominate beam spin asymmetries
- We have an underconstrained set of equations
  - If we want to get to the 3D picture of the nucleon, we will need some tricks

$$\Delta \sigma_{LU} \propto \sin \phi \ Im\{F_1 \mathcal{H} + \xi(F_1 + F_2)\tilde{\mathcal{H}} - kF_2 \mathcal{E} + ...\}$$

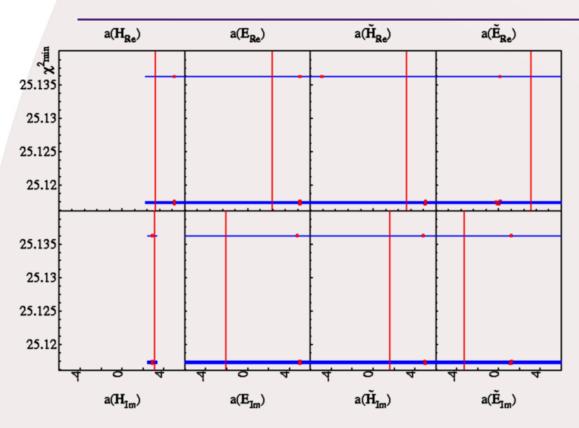
$$\Delta \sigma_{UL} \propto \sin \phi \ Im\{F_1 \tilde{\mathcal{H}} + \xi(F_1 + F_2)\left(\tilde{\mathcal{H}} + \frac{x_B}{2}\mathcal{E}\right) - \xi kF_2 \tilde{\mathcal{E}} + ...\}$$

$$\Delta \sigma_{LL} \propto (A + B\cos\phi) \ Re\{F_1 \tilde{\mathcal{H}} + \xi(F_1 + F_2)\left(\mathcal{H} + \frac{x_B}{2}\mathcal{E}\right) + ...\}$$

$$\Delta \sigma_{Ux} \propto \sin\phi \ Im\{k(F_2 \mathcal{H} - F_1 \mathcal{E}) + ...\}$$



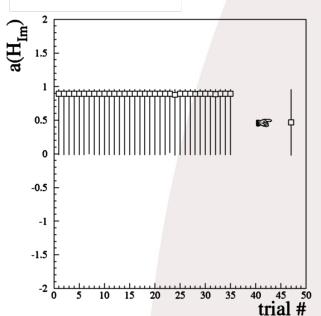
## **Tests with Pseudo Data**

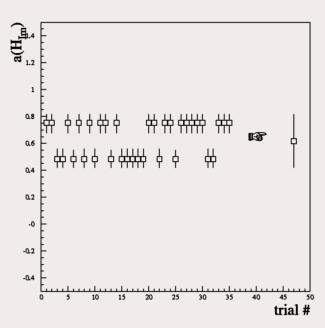


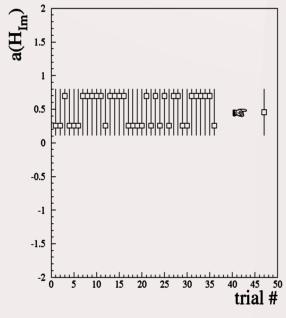
- As expected the 8 parameter fit cannot converge without constraints
  - Choice of educated limits for the local fit  $\rightarrow$  ± 5 x VGG Model predictions
  - As expected H appears to be the only GPD constrained by CLAS data
- Fit sometimes cannot find a single unique solution
  - We ran many fits on pseudo data in order to test the stability of the fitting procedure
    - → Whatever we use in the start, we recover H properly (within error bars) at the end



# **Using central values**





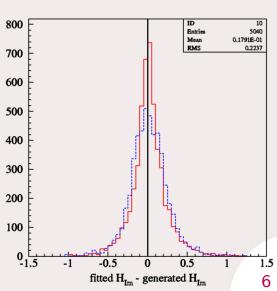


#### In some cases the fit gives problematic results

- Explored with many independent pseudo-data sets and fit starting points
- Highly asymmetric error bars
  - Do not reflect properly the χ<sup>2</sup> profile
  - De to very flat χ<sup>2</sup> valley
- Double solutions
  - Problematic for the coming global fit, which solution to use?

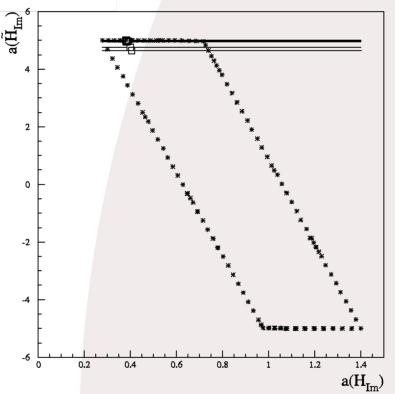
#### We found that the central value of the error bars works best

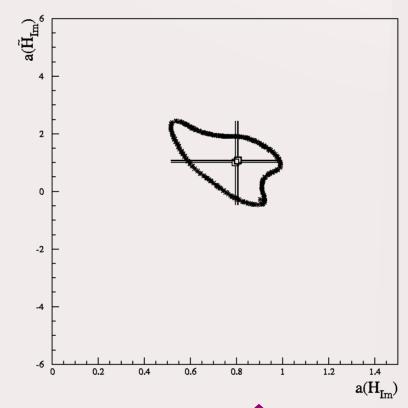
- This is natural since subleading CFFs are in fact not significantly constrained and the minimum  $\chi^2$  in their range is most of the time not significant
- This was confirmed by simulation





# Importance of new observables





- Adding target asymmetries constrains naturally the Im(Ĥ)
  - Incidentally it also constrains Im(H)
- However these data are not available for all kinematics
  - More observables would be needed to constrain E and E
  - Transversely polarized target and charge asymmetries for example

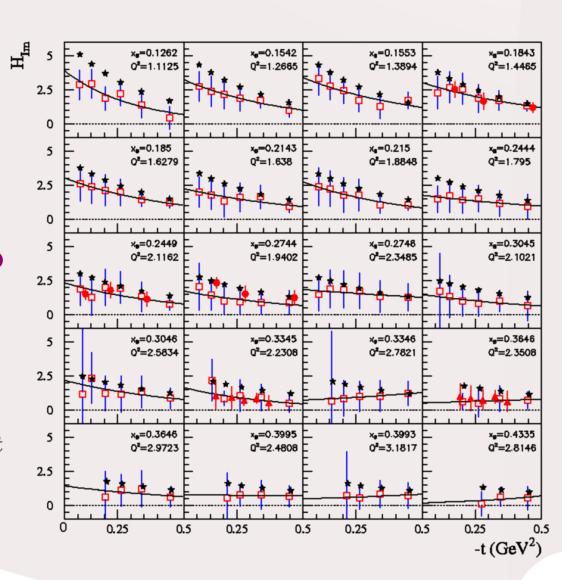


# **Extraction of Im(H)**

- Applying the local fit method for the JLab data
  - Jlab Hall A  $(\sigma, \Delta\sigma)$
  - CLAS  $(\sigma, \Delta \sigma, ITSA, DSA)$
- Gives enough coverage to explore the t and  $x_B (\rightarrow \xi)$  dependence
  - Fitted with an exponential form

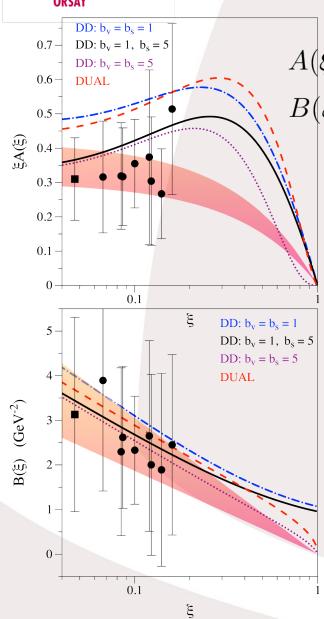
$$\mathcal{H}_{Im}(\xi, t) = A(\xi)e^{B(\xi)t}$$

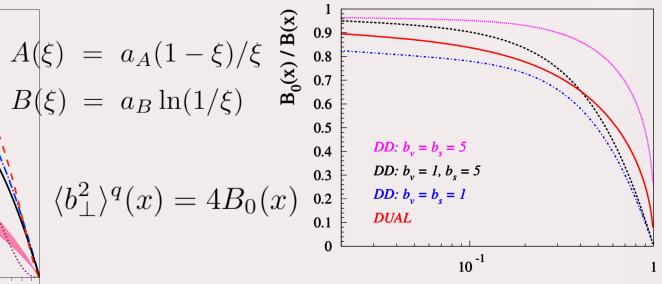
 Results are generally close to the VGG model





# **Amplitude and Slope**





#### The A and B parameters contain the physics

- Linked to density and to transverse size

#### Fitted using educated guess

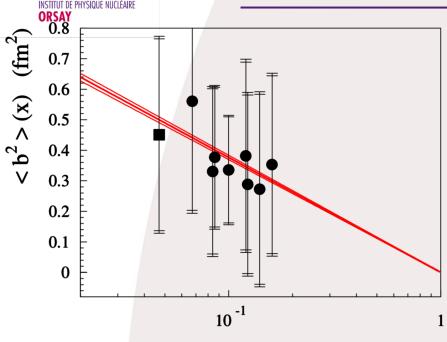
- Asymptotic behavior expectations similar to PDFs
- In the future with larger amount of data, models can be tested at this level

#### • The tomography of the nucleon

- Not directly accessible with DVCS, need a  $\xi$  dependent correction to go from the singlet to the non-singlet distribution
- We note that at low x the correction is small and similarly described by several models



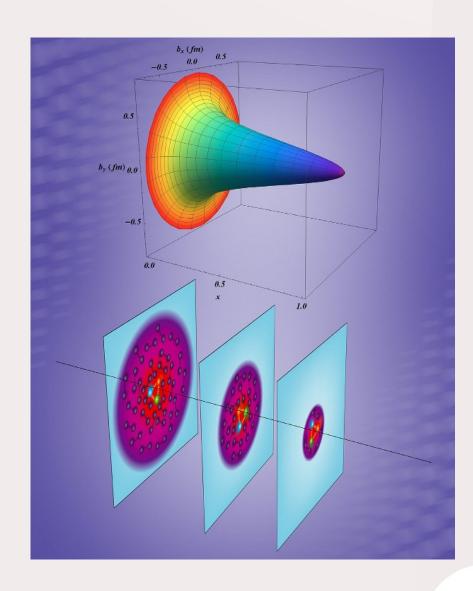
# Results



- We obtain the tomography of the proton
  - Mean square charge radius of the proton for slices of x

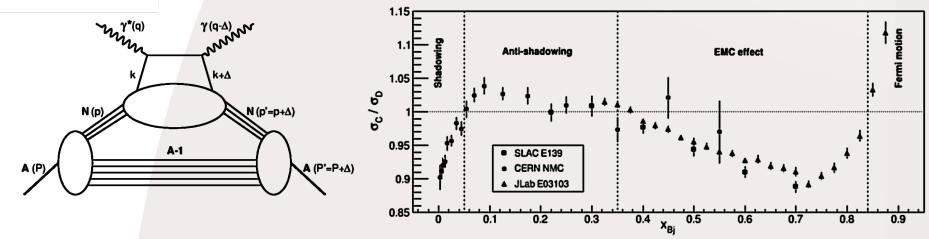
X

- Error bars reflect the unknown CFFs
  - To flatten this distribution, one would need a non constrained CFF with very strong unexpected behavior





## **Nuclear GPDs**



#### New view on nuclear effects

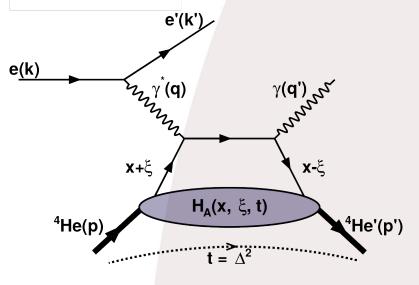
- GPDs offer a completely new point of view into this problem
- Experimental access to completely new nuclear physics
  - Non nucleonic degrees of freedom of the nuclei
  - Measure pressure and forces in the nuclei
  - Localize the EMC effect
     R.D. & S. Scopetta Eur.Phys.J. A52 (2016) no.6, 159

#### Allows to play with the spin

- Use of helium 4 for a simplified problem with only 1 GPD
- Which can be extracted from DVCS beam asymmetries only
- Use of helium 3 and deuterium to understand the neutron and more complex spin dynamics



## **DVCS** on Nuclei

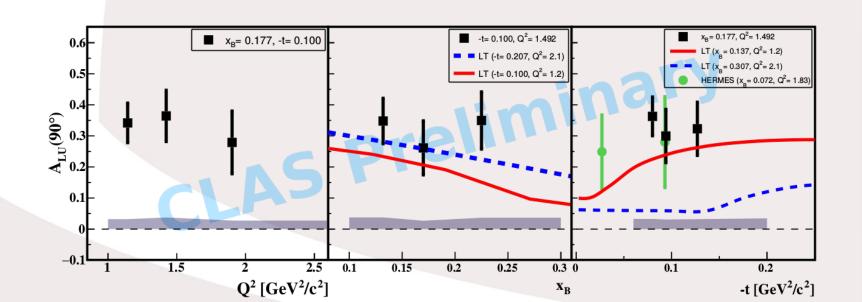


#### Already measured at Jlab (CLAS)

- Coherent DVCS is cleanly measurable
  - Collider kinematics would make it much easier!
- Asymmetries are much larger than for the proton
  - As expected from theory

#### The start of a new domain for GPD studies

- Already many studies on the theory side
  - In both valence and shadowing regions





## **Future at JLab 12**



#### A very wide approved program

- Proton DVCS ( $\sigma$ ,  $\Delta \sigma$ ) in Hall A (running already) and CLAS12
  - Also with longitudinally polarized target
  - Q<sup>2</sup> scan in Hall C, beam energy scan in CLAS12
- Neutron DVCS ( $\sigma$ ,  $\Delta \sigma$ ) in CLAS12
  - Also with longitudinally polarized target
- Transversely polarized target measurement (HD Ice target)
- And many other projects proposed

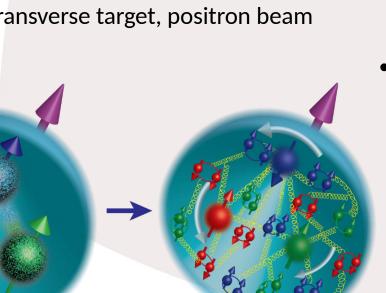
#### JLab will provide a very large data set in the valence region

We should try to equal this for the sea region at EIC



# **Summary**

- We make the tomography of the nucleon
  - Errors can be reduced by including more observables
    - Cross-sections, beam spin asymmetries, target asymmetries...
    - Transverse target, positron beam



- We show for the first time that the proton is smaller at higher x
  - with minimum model input
- What does it says for EIC :
  - We need high luminosity
  - High polarization of both beams
- What to do to go beyond
  - Measure other processes
    - Double DVCS, Time-like CS...
    - Charge asymmetries, transverse polarized target